

Extremal quantum protocols

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Generalized quantum instruments correspond to measurements where the input and output are either states or more generally quantum circuits. These measurements describe any quantum protocol including games, communications, and algorithms. The set of generalized quantum instruments with a given input and output structure is a convex set. Here, we investigate the extremal points of this set for the case of finite dimensional quantum systems and generalized instruments with finitely many outcomes. We derive algebraic necessary and sufficient conditions for extremality. © 2011 American Institute of Physics. [doi:10.1063/1.3610676]

I. INTRODUCTION

Experiments in quantum theory can be modeled through quantum networks that provide the natural description of an arbitrary quantum procedure, corresponding to a causal sequence of steps. The most basic building blocks of quantum networks are state preparations, state transformations (channels and state reductions) and measurements. Provided we have a quantum network, we can isolate open sub-circuits, whose connections constitute the whole network. Any optimization problem in quantum theory can be seen as the search for the most suitable sub-circuit for a specified purpose. For example, for discrimination of states we need to optimize a measurement, or for discrimination of channels we need to optimize the network into which the channel is inserted. Open sub-circuits provide a representation for the most general quantum protocol, where the gates represent the sequence of operations performed by the agent that is communicating, computing or applying a strategy for a quantum game. From a more abstract point of view any sub-circuit represents the most general input-output map that can be achieved via a quantum circuit, that is called generalized quantum instrument (GQI).¹ GQIs then provide the mathematical description for any quantum protocol including games, communications, and algorithms. It is possible to uniquely associate² a positive operator to any deterministic GQI—corresponding to a sub-circuit that does not provide outcomes—in the same way as a positive operator is associated to any channel through the Choi-Jamiołkowski correspondence. More generally, it is possible to associate a set of positive operators to any GQI (Ref. 3) in such a way that each operator corresponds to a possible measurement outcome and summarizes the probabilistic input-output behavior of the GQI as a sub-circuit, conditionally on the outcome. The advantage of this description comes from neglecting the implementation details that are irrelevant for the input-output behavior of the GQI within a quantum network, like arbitrary transformations on ancillary systems, etc. The set of GQIs with the same input and output types is convex, since a random choice of two different GQIs provides a convex combination of the corresponding two input-output maps. It is thus clear that the description of quantum maps through GQIs¹ in optimization problems is convenient for two reasons. The first one is that this approach gets rid of many irrelevant parameters, and the second one is that the optimization problems are reduced to convex optimization on suitably defined convex sets. Applications of GQIs in optimization problems

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can be found in Refs. 4–11. The theory of GQIs was alternatively introduced¹ as a theory of higher order quantum functions, spawning interest in the investigation of more computational consequences of the properties of GQIs.^{12–14} A similar approach to general affine functions on convex subsets of state spaces was recently published,¹⁵ explicitly inspired to the concept of GQIs and quantum combs (namely singleton GQIs).

As a special case of GQIs, we have the elementary examples of states, channels, and positive operator valued measures (POVMs). The analysis of the extremality conditions for states is trivial, and can be found in any textbook of quantum theory. Algebraic extremality conditions for channels were provided in Ref. 16, while the conditions for POVMs were derived later in Refs. 17–21. Other special cases of GQIs are quantum combs,² corresponding to deterministic GQIs, or quantum testers,^{1,3} which are GQIs with outputs that are probability distributions. While all GQIs could be decomposed into states, channels, and measurements, it is much more practical to consider the corresponding networks as a whole.

Optimization tasks in quantum information processing can be rephrased in terms of optimization of a certain GQI with respect to some particular figure of merit, which is often a convex function on the set of GQIs and the maximum is achieved on an extremal point of this set. Moreover, also for those problems that resort to convex optimization or minimax problems, numerical optimization is enhanced by the possibility of generating arbitrary extremal elements. For this purpose, having an algebraic characterization is a crucial step.

In the present paper, we consider the convex sets of GQIs, and characterize their extremal points for the case of finite dimensional quantum systems and the instruments that have finitely many outcomes. As special cases we obtain the extremality conditions for POVMs, channels, testers or instruments.

The paper is organized as follows. In Sec. II, we introduce the theoretical framework we use to describe quantum networks. In Sec. III, we formulate the necessary and sufficient condition of extremality for GQI. Sections IV, V, and VI study the implications of the extremality condition in the case of quantum testers, quantum channels, and quantum instruments, respectively. Finally, the summary of the results is placed in Sec. VII.

II. THEORY OF QUANTUM NETWORKS

Let us summarize some pieces of the theoretical framework of quantum networks introduced in Ref. 1 that we will use. An arbitrary quantum network \mathcal{R} can be formally understood as a quantum memory channel,² whose inputs and outputs are labeled by even or odd numbers from 0 to $2N - 1$, respectively. The Hilbert spaces associated with these inputs and outputs can be in general different and we denote them by \mathcal{H}_i $i = 0, \dots, 2N - 1$. As it was shown already in Ref. 2, deterministic quantum network \mathcal{R} is fully characterized by its Choi-Jamiołkowski operator, i.e., a deterministic quantum N -comb R .

Definition 1: A deterministic quantum N -comb on $\mathcal{H}_0, \dots, \mathcal{H}_{2N-1}$ is a positive operator $R \equiv R^{(N)} \in \mathcal{L}(\mathcal{H}_0 \otimes \dots \otimes \mathcal{H}_{2N-1})$, which obeys the following normalization conditions:

$$\begin{aligned} \text{Tr}_{2n-1} R^{(n)} &= I_{2n-2} \otimes R^{(n-1)}, \quad 0 \leq n \leq N, \\ \text{Tr}_1 R^{(1)} &= I_0, \end{aligned} \quad (1)$$

where the operators $R^{(n)}$ are defined recursively.

Positive operators $T \in \mathcal{L}(\mathcal{H}_0 \otimes \dots \otimes \mathcal{H}_{2N-1})$, such that $T \leq R$ for some deterministic quantum comb R , are called non-deterministic quantum N -combs. An arbitrary probabilistic quantum network, whose different outcomes are indexed by $i = 1, \dots, M$ is described by a collection of non-deterministic quantum N -combs $\{T_i\}_{i=1}^M$ defined as follows:

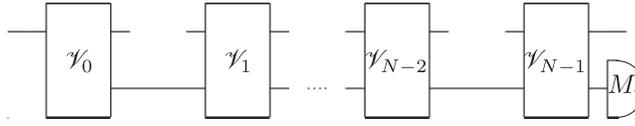


FIG. 1. The circuit implementation of a general GQI shows that GQIs correspond to the most general quantum network,¹ and vice versa one can prove that any quantum circuit possibly involving measurements corresponds to a GQI. The transformations \mathcal{V}_i are isometries, and M_i denotes a POVM. Notice that the measurement can always be postponed to the very last step.

Definition 2: Generalized quantum N -instrument is a collection $\{T_i\}_{i=1}^M$ of non-deterministic quantum N -combs that sum up to a deterministic quantum comb,

$$\sum_{i=1}^M T_i = R. \quad (2)$$

A realization theorem can be proved,¹ providing the interpretation of GQIs as the appropriate mathematical representation for the most general quantum network, because any GQI can be implemented through a quantum circuit as in Fig. 1, and vice versa any quantum circuit possibly involving measurements corresponds to a GQI.

For $M = 1$ the corresponding network is deterministic and the set of generalized quantum instruments coincides with deterministic quantum combs. On the other hand, if $N = 1$ a generalized quantum instrument is a collection of completely positive maps forming a channel, which is usually called an instrument. Another special case of generalized quantum instruments is provided by quantum testers.

Definition 3: A quantum N -tester is a generalized quantum $(N + 1)$ -instrument on $\mathcal{H}_0, \dots, \mathcal{H}_{2N+1}$ with one-dimensional Hilbert spaces $\mathcal{H}_0, \mathcal{H}_{2N+1}$.

Quantum testers are analogous to the concept of positive operator valued measures (POVM) as they allow to express probability distributions for arbitrary tests on quantum combs.

We will show our analysis of the extremal points of the set of generalized quantum instruments, which provides necessary and sufficient conditions for extremality and leads to specific new conditions also for all the above mentioned special cases.

III. EXTREMALITY CONDITION FOR GENERALIZED QUANTUM INSTRUMENTS

In this section, we shall apply the method of perturbations to find extremal generalized quantum instruments. The perturbation method was also used to determine extremal channels¹⁶ and POVMs.¹⁹ However, the application of the perturbation method to GQIs does not come as a straightforward generalization of previous results, because the richer structure of the normalization constraints for GQIs requires a radically different analysis.

Let us consider arbitrary generalized quantum N -instrument $\{T_i\}_{i=1}^M$. We denote by \mathcal{V}_i the support of the operator T_i . The support of the sum of positive operators is the span of the supports of the summed operators. Thus, the support of the normalization $R = \sum_{i=1}^M T_i$ is $\mathcal{H}_R \equiv \text{Span}\{\mathcal{V}_i\}_{i=1}^M$. A set of operators $\{D_i\}_{i=1}^M$ is called a valid perturbation of GQI $\{T_i\}_{i=1}^M$ if and only if $\{T_i \pm D_i\}_{i=1}^M$ are valid GQIs. Existence of a perturbation has two major implications. First, the positivity of $T_i \pm D_i$ requires D_i to be hermitian and to have support only in \mathcal{V}_i . This is proved by the following lemma.

Lemma 1: Suppose that operators T, D fulfill $T \geq 0, D^\dagger = D$. If $T \geq \pm D$, then $\text{Supp}(D) \subseteq \text{Supp}(T)$.

Proof: The statement of the lemma can be equivalently formulated as $\text{Ker}(T) \subseteq \text{Ker}(D)$. This can be proved considering the general decomposition of a vector $|\psi\rangle$ as $\alpha|\Psi_S\rangle + \beta|\Psi_K\rangle$ where $\Psi_S \in \text{Supp}(T)$ and $\Psi_K \in \text{Ker}(T)$. Then we have

$$|\alpha|^2 \langle \Psi_S | (T \pm D) | \Psi_S \rangle \pm 2\Re \alpha^* \beta \langle \Psi_S | D | \Psi_K \rangle \pm |\beta|^2 \langle \Psi_K | D | \Psi_K \rangle \geq 0, \quad (3)$$

for all α, β . Choosing $\alpha = 0$ one immediately obtains $\langle \Psi_K | D | \Psi_K \rangle = 0$, which by the polarization identity implies also $\langle \Psi'_K | D | \Psi_K \rangle = 0$ for all $\Psi'_K \in \text{Ker}(T)$. The previous inequality can thus be rewritten as follows

$$|\alpha|^2 \langle \Psi_S | (T \pm D) | \Psi_S \rangle \pm 2\Re \alpha^* \beta \langle \Psi_S | D | \Psi_K \rangle \geq 0, \quad (4)$$

for all α, β . Suitably choosing the phases of α, β , one has

$$|\alpha|^2 \langle \Psi_S | (T \pm D) | \Psi_S \rangle \mp 2|\alpha||\beta| |\langle \Psi_S | D | \Psi_K \rangle| \geq 0, \quad (5)$$

and for $\beta = \frac{1}{2}$ and $|\alpha| > 0$ we obtain

$$|\alpha| |\langle \Psi_S | (T \pm D) | \Psi_S \rangle| \geq |\langle \Psi_S | D | \Psi_K \rangle|, \quad (6)$$

for all $|\alpha|$, implying that $\langle \Psi_S | D | \Psi_K \rangle = 0$ holds for all Ψ_S and Ψ_K . This together with $\langle \Psi'_K | D | \Psi_K \rangle = 0 \forall \Psi'_K \in \text{Ker}(T)$ allows us to conclude that $\langle \psi | D | \Psi_K \rangle = 0$ for every $\psi \in \mathcal{H}$, i.e., $D | \Psi_K \rangle = 0$ for all Ψ_K . This proves that $\text{Ker}(T) \subseteq \text{Ker}(D)$, or, equivalently, $\text{Supp}(D) \subseteq \text{Supp}(T)$. ■

As a consequence if we write operators T_i in their spectral form $T_i = \sum_k \lambda_k^{(i)} |v_k^i\rangle \langle v_k^i|$ then arbitrary hermitian operator D_i with support in \mathcal{V}_i can be written as

$$D_i = \sum_{n,m} D_{nm}^{(i)} |v_n^i\rangle \langle v_m^i|, \quad (7)$$

where $D_{nm}^{(i)}$ is a hermitian matrix with $r_i^2 \equiv (\dim \mathcal{V}_i)^2$ real parameters. We form a basis $\mathbb{H}_i \equiv \{Q_j^{(i)}\}_{j=1}^{r_i^2}$ of hermitian operators with support in \mathcal{V}_i and we define $\mathbb{D}_M := \bigcup_{i=1}^M \mathbb{H}_i$.

The second consequence of requiring valid perturbed GQI $\{T_i \pm D_i\}_{i=1}^M$ is that, due to the normalization condition (2) the perturbed GQI has to sum up to deterministic N -combs R_{\pm} , which can be stated as

$$\sum_{i=1}^M D_i = \Delta, \quad (8)$$

where $\Delta \equiv \pm(R_{\pm} - R)$ is an operator expressible as a difference of two deterministic quantum N -combs. Using the parametrization of deterministic quantum combs developed in Appendix A it is clear that Δ lies in \mathcal{W}_C , the subspace of operators spanned by the basis

$$\mathbb{D}_{(N)} \equiv \{E_i^{(2N-1)} \otimes F_j^{(2N-2)}, I_{2N-1, 2N-2} \otimes E_i^{(2N-3)} \otimes F_j^{(2N-4)}, \dots, I_{2N-1, \dots, 2} \otimes E_i^{(1)} \otimes F_j^{(0)}\}$$

where $\{E_i^{(k)}\}_{i=2}^{d_k^2}$ is a basis of traceless hermitian operators on \mathcal{H}_k , and $\{F_j^{(k)}\}_{j=1}^{d_k^2 \dots d_0^2}$ is basis of all hermitian operators acting on $\mathcal{H}_k \otimes \mathcal{H}_{k-1} \otimes \dots \otimes \mathcal{H}_0$. On the other hand, due to the positivity requirement for $R_{\pm} = R \pm \Delta$, Δ must be a hermitian operator with support in \mathcal{H}_R . Let us call \mathcal{W}_S the subspace of hermitian operators with support in \mathcal{H}_R . Thus, the allowed perturbations of the normalization lie in the intersection $\mathcal{W}_I \equiv \mathcal{W}_S \cap \mathcal{W}_C$. The relation between non-existence of a valid perturbation and the requirements on the operators D_1, \dots, D_M, Δ is expressed by the following theorem.

Theorem 1: A generalized quantum N -instrument $\{T_i\}_{i=1}^M$ acting on $\mathcal{H}_{2N-1} \otimes \dots \otimes \mathcal{H}_0$ is extremal if and only if $\mathbb{D}_M \cup \mathbb{D}_{(N)}$ is a linearly independent set of operators.

Proof: We are going to prove the theorem by showing the equivalence of the negated statements, i.e., a GQI is not extremal if and only if the basis $\mathbb{D}_M \cup \mathbb{D}_{(N)}$ is linearly dependent. It is easy to show that if a GQI is not extremal, then the basis $\mathbb{D}_M \cup \mathbb{D}_{(N)}$ is linearly dependent. If a point of a convex set is not extremal then there exists a bidirectional perturbation to it. Hence, there exists a set of operators D_i such that $\{T_i \pm D_i\}_{i=1}^M$ is a valid GQI. In particular, due to at least one operator D_i being non-zero we have Eq. (8), which after expanding the LHS in \mathbb{D}_M and RHS in $\mathbb{D}_{(N)}$ proves the linear dependence of basis $\mathbb{D}_M \cup \mathbb{D}_{(N)}$.

In order to prove the converse statement we are going to show that if the basis $\mathbb{D}_M \cup \mathbb{D}_{(N)}$ is linearly dependent then there exists a valid perturbation of the considered GQI and hence it is not

extremal. Linear dependence of $\mathbb{D}_M \cup \mathbb{D}_{(N)}$ means there exists a non-zero vector consisting of all coefficients $D_{nm}^{(i)}$, s_k such that

$$\sum_{i,n,m} D_{nm}^{(i)} |v_n^i\rangle \langle v_m^i| + \sum_k s_k G_k = 0, \quad (9)$$

where $D_{nm}^{(i)}$ are for each i hermitian matrices, $|v_n^i\rangle$ are eigenvectors of T_i and G_k are basis elements of $\mathbb{D}_{(N)}$. Let us recall that the basis $\mathbb{D}_{(N)}$ is by construction linearly independent, so all $D_{nm}^{(i)}$ can not be zero simultaneously. We rewrite the equation (9) as

$$\sum_{i,n,m} D_{nm}^{(i)} |v_n^i\rangle \langle v_m^i| = \sum_k -s_k G_k \equiv \Delta. \quad (10)$$

For each i the operators on the LHS of (10) have support in the subspace \mathcal{V}_i . All subspaces \mathcal{V}_i are included in the support of the normalization R . Thus, the operator on the LHS of (10) belongs to an operator subspace \mathcal{W}_S . Since the RHS of (10) is from subspace \mathcal{W}_C it is clear that $\sum_k s_k G_k \in \mathcal{W}_S \cap \mathcal{W}_C = \mathcal{W}_I$. This implies that for suitably small ε the operator $R \pm \varepsilon \Delta$ is positive as well as all operators $T_i \pm \varepsilon D_i$. Thus, we have found a valid perturbation of the GQI $\{T_i\}_{i=1}^M$ showing that it is not extremal, which concludes the proof. ■

IV. EXTREMALITY OF QUANTUM TESTERS

In this section, we focus our attention to quantum testers, which can be used to solve problems like discrimination of quantum channels, or optimization of quantum oracle calling algorithms and others, because they describe achievable probability distributions for all possible experiments with given resources. More precisely, we consider quantum N -testers and we try to identify the extremal points of this set. We start by the analysis of 1-testers, also called Process-POVMs.²² A 1-tester with M outcomes is defined by positive operators $\{T_i\}_{i=1}^M$ acting on $\mathcal{H}_2 \otimes \mathcal{H}_1$, which satisfy the normalization condition

$$\sum_{i=1}^M T_i = I_2 \otimes \rho_1, \quad (11)$$

where ρ is a state on \mathcal{H}_1 .²³ As before we denote by \mathcal{V}_i the supports of operators T_i . Let us denote the support of ρ by \mathcal{H}_ρ and by $r = \dim \mathcal{H}_\rho$ the rank of ρ . The 1-tester $\{T_i\}_{i=1}^M$ on $\mathcal{H}_2 \otimes \mathcal{H}_1$ can be considered as a valid 1-tester on $\mathcal{H}_2 \otimes \mathcal{H}'_1$ for arbitrary \mathcal{H}'_1 that includes \mathcal{H}_ρ (e.g., $\mathcal{H}'_1 = \mathcal{H}_\rho$).

A. Extremality condition for 1-testers

In the following, we express the general extremality condition from Theorem 1 for 1-testers and we propose a slightly different extremality condition, which is easier to check. The set $\mathbb{D}_{(N)}$ from Theorem 1 is in this case formed by the operators $\{I_2 \otimes E_i^{(1)}\}_{i=2}^{d_1^2}$, where $\{E_i^{(1)}\}_{i=2}^{d_1^2}$ is a basis of trace zero hermitian operators on \mathcal{H}_1 .

Corollary 1: A quantum 1-tester $\{T_i\}_{i=1}^M$ is extremal if and only if there exists only a trivial solution of an equation $\sum_{i=1}^M \sum_{n,m} D_{nm}^{(i)} |v_n^i\rangle \langle v_m^i| + \sum_{j=1}^{d_1^2-1} s_j I_2 \otimes E_j^{(1)} = 0$, where $\forall i$ $D_{nm}^{(i)}$ are hermitian matrices and s_j are real numbers.

Since the normalization of the perturbed tester must be supported inside the support of the original normalization, it is natural that, $\mathbb{D}_I \equiv \{I_2 \otimes \sigma_l\}_{l=1}^{r^2-1}$, the basis of trace zero operators supported under the original normalization $I_2 \otimes \rho_1$ can be used in the Theorem 1 instead of $\mathbb{D}_{(N)}$.

Theorem 2: A quantum 1-tester $\{T_i\}_{i=1}^M$ is extremal if and only if the equation

$$\sum_{i=1}^M \sum_{n,m} D_{nm}^{(i)} |v_n^i\rangle \langle v_m^i| + \sum_{l=1}^{r^2-1} s_l I_2 \otimes \sigma_l = 0, \quad (12)$$

where $D_{nm}^{(i)}$ are for each i hermitian matrices and s_i are real numbers, has only a trivial solution.

Actually, the basis \mathbb{D}_I of the subspace \mathcal{W}_I can be always used in the Theorem 1 and the proof still holds. However, for $N \neq 1, 2$ it is often easier to specify $\mathbb{D}_{(N)}$ rather than \mathbb{D}_I .

As we said for 1-testers \mathbb{D}_I is formed by trace zero operators supported under ρ tensored with unity on \mathcal{H}_2 and this will help us to get more insight to 1-testers. The extremality condition for 1-tester $\{T_i\}_{i=1}^M$ from Theorem 2 allows us to give the following bound

$$\sum_{i=1}^M r_i^2 + r^2 - 1 \leq (rd_2)^2 \tag{13}$$

on the ranks r_i of the operators T_i . The bound is derived by counting the number of elements of $\mathbb{D}_M \cup \mathbb{D}_I$ and realizing that these operators should be linearly independent hermitian operators acting only on $\mathcal{H}_2 \otimes \mathcal{H}_\rho$. From the bound (13) it is clear that the extremal tester can have the highest possible number of outcomes if $r = d_1$ and the ranks r_i are as close to one as possible. Assuming all r_i are rank 1 we get the bound on the number of elements of the extremal quantum 1-tester

$$M \leq d_1^2(d_2^2 - 1) + 1, \tag{14}$$

B. Classification of extremal 1-testers

Let us now answer the question, which normalizations $I \otimes \rho$ allow existence of extremal testers. For this purpose let us define a superoperator $\xi_{\rho,U}$ that acts on linear operators on $\mathcal{H}_2 \otimes \mathcal{H}_1$ as

$$\xi_{\rho,U}(T_i) \equiv d_1(I \otimes \sqrt{\rho} U) T_i (I \otimes U^\dagger \sqrt{\rho}). \tag{15}$$

For any state ρ with full rank (i.e., $r = d_1$) and any unitary U acting on \mathcal{H}_1 , the superoperator $\xi_{\rho,U}$ is invertible and preserves positivity of operators. Using $\xi_{\rho,U}$ we can formulate the following theorem.

Theorem 3: Suppose we have a full rank state ρ , a unitary operator U and a 1-tester $\{T_i\}_{i=1}^M$ on $\mathcal{H}_2 \otimes \mathcal{H}_1$ with $\sum_{i=1}^M T_i = I \otimes \frac{1}{d_1} I$. Then the tester $\{T'_i \equiv \xi_{\rho,U}(T_i)\}_{i=1}^M$ on $\mathcal{H}_2 \otimes \mathcal{H}_1$ has normalization $\sum_{i=1}^M T'_i = I \otimes \rho$ and is extremal if and only if 1-tester $\{T_i\}_{i=1}^M$ is extremal.

Proof: First, let us note that the form of $\xi_{\rho,U}$ guarantees positivity of T'_i and leads to the normalization

$$\sum_{i=1}^M T'_i = \xi_{\rho,U}(I \otimes \frac{1}{d_1} I) = I \otimes \rho. \tag{16}$$

Now we prove that the tester $\{T'_i\}_{i=1}^M$ is extremal if the original tester $\{T_i\}_{i=1}^M$ was. Let us stress that for any extremal 1-tester its normalization $I \otimes \rho$ is (up to multiplication) the only operator of the form $I \otimes X$ that is in the span of the operators D_i . This holds, because the span of the operators $D_i \in \mathcal{L}(\mathcal{V}_i) \subseteq \mathcal{L}(\mathcal{H}_2 \otimes \mathcal{H}_\rho)$ covers $I \otimes \rho$ and it is independent from $r^2 - 1$ dimensional subspace of traceless hermitian operators of the form $I \otimes X$ due to linear independence (12). Superoperator $\xi_{\rho,U}$ is invertible so it preserves linear independence. In our case this means that the basis \mathbb{H}'_i of hermitian operators derived from $\xi_{\rho,U}(|v_n^i\rangle\langle v_m^i|)$ is linearly independent and spans the whole space of hermitian operators that have support in the support of T'_i . Moreover, due to extremality of the original tester $\{T_i\}_{i=1}^M$ ($\sum_{i=1}^M T_i = I \otimes \frac{1}{d_1} I$) and the invertibility of $\xi_{\rho,U}$ we can conclude that also $I \otimes \rho = \xi_{\rho,U}(I \otimes \frac{1}{d_1} I)$ is the only operator of the form $I \otimes X$ that is in the span of $\mathbb{D}'_M = \cup_{i=1}^M \mathbb{H}'_i$. Let us assume that the tester $\{T'_i\}_{i=1}^M$ is not extremal even though the original tester $\{T_i\}_{i=1}^M$ was extremal. In other words we assume that \mathbb{D}'_M is linearly dependent with traceless operators $\{I_2 \otimes \sigma_i\}_{i=1}^{r^2-1}$. As a consequence there must exist a traceless operator of the form $I \otimes X$ in the span D'_i . However, this is a contradiction, because the only operator of such form is $I \otimes \rho$ and has trace one. We conclude that the transformed tester $\{T'_i\}_{i=1}^M$ must be extremal.

In fact, the same argumentation can be used to prove that $\{T_i\}_{i=1}^M$ is extremal if $\{T'_i\}_{i=1}^M$ was, because $\xi_{\rho,U}$ is invertible. Hence, for arbitrary extremal tester $\{T'_i\}_{i=1}^M$ with normalization $I \otimes \rho$ using $(\xi_{\rho,U})^{-1}$ one obtains extremal tester $\{T_i\}_{i=1}^M$ with normalization $I \otimes \frac{1}{d_1}I$. ■

The Theorem 3 is very useful, because to classify all extremal 1-testers it suffices to classify extremal 1-testers with normalization $I \otimes \frac{1}{d_1}I$. More precisely, using $\xi_{\rho,I}^{-1}$ each extremal tester is in one to one correspondence with an extremal tester with normalization $I \otimes \frac{1}{r}\Pi_\rho$, where Π_ρ is a projector onto a support of ρ . This tester can be considered as a tester on $\mathcal{H}_2 \otimes \mathcal{H}_\rho$, where its normalization is of the above mentioned form $I \otimes \frac{1}{d_1}I$. Thus, we can formulate the following corollary of Theorem 3.

Corollary 2: Extremal 1-testers with M outcomes exist either for all normalizations $I \otimes \rho$ with given rank r of ρ or for none of them.

Let us now relate the set $\Theta(\mathcal{H}_2, \mathcal{H}_1)$ of extremal quantum testers with normalization $I \otimes \frac{1}{d_1}I$ to the set $P(\mathcal{H}_2 \otimes \mathcal{H}_1)$ of extremal POVMs on $\mathcal{H}_2 \otimes \mathcal{H}_1$. Namely, each extremal tester $\{T_i\}_{i=1}^M$ with normalization $I \otimes \frac{1}{d_1}I$ defines an extremal POVM $\{E_i = d_1 T_i\}_{i=1}^M$. This follows directly from the extremality condition for quantum testers (12), which necessarily requires the basis of hermitian operators with supports on \mathcal{V}_i to be linearly independent. This is exactly the necessary and sufficient condition for the extremality of the POVM (Ref. 19) $\{E_i\}_{i=1}^M$. Apart from the multiplicative difference in normalization, we will prove later that extremal quantum testers with normalization $I \otimes \frac{1}{d_1}I$ are a proper subset of extremal POVMs on $\mathcal{H}_2 \otimes \mathcal{H}_1$. On the other hand there are extremal POVMs on $\mathcal{H}_2 \otimes \mathcal{H}_1$, which cannot be rescaled to form an extremal tester. One example are informationally complete POVMs on $\mathcal{H}_2 \otimes \mathcal{H}_1$ with $(d_1 d_2)^2$ outcomes. Their existence was proved in Ref. 19 for any dimension, but they have too many outcomes to form an extremal 1-tester (see Eq. (14)).

C. Extremal 1-testers with rank one normalization

Having a tester with rank one normalization $\rho = |\phi\rangle\langle\phi|$ implies that all the elements of the tester have the form $T_i = E_i \otimes \rho$, where E_i is positive operator acting on \mathcal{H}_2 . Let us note that these testers correspond to preparation of a pure state ρ and performing a POVM $\{E_i\}_{i=1}^M$. Since the support of ρ is one-dimensional, there are no traceless operators with support in \mathcal{H}_ρ . Thus, the extremality condition (12) is in this case equivalent to linear independence requirement

$$0 = \sum_{i=1}^M \sum_{n,m} D_{nm}^{(i)} |w_n^i\rangle\langle\phi| \langle w_m^i| \langle\phi| = \left(\sum_{i=1}^M \sum_{n,m} D_{nm}^{(i)} |w_n^i\rangle\langle w_m^i| \right) \otimes |\phi\rangle\langle\phi| \Rightarrow D_{nm}^{(i)} = 0 \quad \forall i, n, m$$

for the basis of hermitian operators on the supports of E_i . This is precisely the necessary and sufficient condition of the extremality of the POVM (Ref. 19) with elements E_i . Thus, the quantum tester $\{T_i = |\phi\rangle\langle\phi| \otimes E_i\}_{i=1}^M$ is extremal if and only if POVM $\{E_i\}_{i=1}^M$ is extremal. In particular, the number of outcomes of the extremal quantum tester in this case cannot exceed d_2^2 , which is the number given by the bound (14) and by the maximal number of elements of an extremal POVM (Ref. 19) as well. On the other hand, a single outcome extremal POVM $\{E_1 = I\}$ leads to an extremal 1-tester $\{T_1 = I \otimes \rho\}$ for arbitrary pure state normalization ρ .

Remark 1: Actually, the only extremal single outcome 1-testers are those with pure state normalization.

D. Extremal qubit 1-testers

For qubit tester ($d_1 = d_2 = 2$) the rank r of the normalization ρ can be either one or two. If ρ is a pure state ($r = 1$) then Sec. IV C tells us that such extremal testers are in one to one correspondence with the extremal qubit POVMs, which can have at most four outcomes. Hence, to classify all extremal qubit testers (based on Sec. IV B) it remains to investigate qubit testers with

normalization $\rho = I \otimes \frac{1}{2}I$. We will identify extremal testers with two outcomes. Then we discuss the case $2 < k \leq 13$ (see bound (14)) and we propose some ways how to construct such testers.

1. Two outcome testers

Considering the ranks r_1, r_2 of the two parts of the tester, there are only three possibilities compatible with bound (13): *i*) $(r_1, r_2) = (1, 3)$, *ii*) $(r_1, r_2) = (2, 2)$, *iii*) $(r_1, r_2) = (2, 3)$, where we assume without loss of generality that $r_1 \leq r_2$. As we already mentioned the supports of the tester operators T_i necessarily have to obey conditions for extremal POVMs on $\mathcal{H}_2 \otimes \mathcal{H}_1$. In particular, operators T_i cannot have intersecting supports (see corollary 3 in Ref 19). This rules out $(r_1, r_2) = (2, 3)$ case.

Let us now consider the case *i*) $(r_1, r_2) = (1, 3)$. In this case T_1 necessarily equals $\frac{1}{2}$ projector onto a pure state, because otherwise the rank of $T_2 = \frac{1}{2}I \otimes I - T_1$ would not be three. Consequently, we can write the tester as

$$T_1 = \frac{1}{2}|\phi_1\rangle\langle\phi_1|, \quad (17)$$

$$T_2 = \frac{1}{2}(I \otimes I - |\phi_1\rangle\langle\phi_1|) = \frac{1}{2} \sum_{i=2}^4 |\phi_i\rangle\langle\phi_i|,$$

where vectors $|\phi_i\rangle$ $i = 1, \dots, 4$ form an orthonormal basis of $\mathcal{H}_2 \otimes \mathcal{H}_1$. As we show in the Appendix B 1, the only two-outcome testers of the above form that are not extremal are those with $|\phi_1\rangle$ being a product state. Looking on how the considered type of testers transforms under superoperator $\xi_{\rho, I}$ from equation (15) one can easily conclude that also for arbitrary rank two normalization ρ the two outcome testers with $(r_1, r_2) = (1, 3)$ are extremal if and only if $|\phi_1\rangle$ is not a product state.

The case *ii*) $(r_1, r_2) = (2, 2)$ has some similarities to the previous one. Since T_1, T_2 are both rank two and their sum is $\frac{1}{2}I \otimes I$, then they both must be equal to $\frac{1}{2}P_i$, where P_i are orthogonal projectors. Consequently, we can write the tester as

$$\begin{aligned} T_1 &= \frac{1}{2}P_1 = \frac{1}{2}(|\phi_1\rangle\langle\phi_1| + |\phi_2\rangle\langle\phi_2|), \\ T_2 &= \frac{1}{2}P_2 = \frac{1}{2}(|\phi_3\rangle\langle\phi_3| + |\phi_4\rangle\langle\phi_4|), \end{aligned} \quad (18)$$

where vectors $|\phi_i\rangle$ $i = 1, \dots, 4$ form an orthonormal basis of $\mathcal{H}_2 \otimes \mathcal{H}_1$. As we show in the Appendix B 2 this type of tester is not extremal only if $P_1 = I \otimes |v\rangle\langle v|$ for some $|v\rangle \in \mathcal{H}_1$ or if the states $|\phi_1\rangle, |\phi_3\rangle$ can be chosen as $|\phi_1\rangle = |w\rangle \otimes |v\rangle, |\phi_3\rangle = |w^\perp\rangle \otimes |v\rangle$ for some states $|w\rangle \in \mathcal{H}_2, |v\rangle \in \mathcal{H}_1$. For arbitrary rank two normalization ρ the conditions on extremality of this type of tester are very similar, but with P_1, P_2 playing the role of projectors onto the support of T_1, T_2 .

2. M-outcome testers

The analysis of extremal qubit testers for more than two outcomes is very involved. For this reason, we provide only some examples how one can construct them. Extremal qubit 1-testers with 3 or 4 outcomes can be easily obtained by taking the extremal 2-outcome tester from Eq. (18) and splitting either one or both its parts into rank one operators. Obviously this operation reduces the subspace achievable by linear combination of operators with support on T_i , thus the linear independence with the operators $\sigma_i \otimes I$ remains untouched and the tester obtained in this way is extremal. A different approach allows us to generate examples of extremal testers with up to $M \leq 10$ as follows. Let us consider an extremal 2-outcome tester from Eq. (17) and let us split its element T_2 into T'_2, \dots, T'_M in such a way that $\{2T'_i\}_{i=2}^M$ is an extremal POVM on the support of T_2 . By setting $T'_1 = T_1$ we obtain an extremal tester $\{T'_i\}_{i=1}^M$, because we are only restricting the operator span of allowed perturbations of the elements T_i and perturbations of T'_2, \dots, T'_M are independent by construction. Finally, one can use the technique of Heinossari and Pellonpää²⁴ (see Proposition 4) to construct extremal qubit 1-testers with $4 \leq M \leq 13$ rank 1 elements. The construction generates

$M + 1$ outcome tester from the M outcome tester until the linear independence of rank 1 elements with the operators $\sigma_i \otimes I$ can be kept (i.e., $M \leq 13$).

V. EXTREMALITY OF QUANTUM CHANNELS

The aim of this section is to show how our general criterion from Theorem 1 in the case of channels ($N = 1$, $M = 1$) relates to known conditions of extremality. For channels mapping from $\mathcal{L}(\mathcal{H}_0)$ to $\mathcal{L}(\mathcal{H}_1)$, we have $\mathbb{D}_{(N)} = \{\sigma_a \otimes I, \sigma_a \otimes \sigma_b\}$, where $\{\sigma_a\}_{a=2}^{d_1^2}$, $\{\sigma_b\}_{b=2}^{d_0^2}$ are basis of trace zero hermitian operators on \mathcal{H}_1 , \mathcal{H}_0 , respectively. Suppose we want to test whether a channel \mathcal{E} with Choi-Jamiolkowski operator E is extremal. If we take the spectral decomposition of $E = \sum_m |K_m\rangle\rangle\langle\langle K_m|$ then the eigenvectors $|K_m\rangle\rangle$ correspond through isomorphism²⁵ $|A\rangle\rangle = A \otimes I|I\rangle\rangle$ (here $|I\rangle\rangle \equiv \sum_i |i\rangle \otimes |i\rangle \in \mathcal{H}_0^{\otimes 2}$) to Kraus operators K_m of a minimal Kraus representation of channel \mathcal{E} . The well-known Choi extremality condition¹⁶ writes

$$\sum_{m,n} \alpha_{mn} K_m^\dagger K_n = 0 \quad \Leftrightarrow \quad \alpha_{mn} = 0 \quad \forall m, n. \quad (19)$$

On the other hand, according to our Theorem 1 the condition for extremality of channel \mathcal{E} is that

$$\begin{aligned} \sum_{m,n} \alpha_{mn} |K_m\rangle\rangle\langle\langle K_n| + \sum_a \beta_a \sigma_a \otimes I + \sum_{a,b} \gamma_{ab} \sigma_a \otimes \sigma_b = 0, \\ \Leftrightarrow \alpha_{mn} = 0, \quad \forall m, n, \quad \beta_a = 0 \quad \forall a, \quad \gamma_{ab} = 0 \quad \forall a, b. \end{aligned} \quad (20)$$

We will now prove the following theorem:

Theorem 4: *The conditions in Eq. (19) and Eq. (20) are equivalent.*

Proof: In order to prove that the condition in Eq. (19) implies the condition in Eq. (20), it is sufficient to suppose that Eq. (20) holds, and to take the partial trace on the Hilbert space \mathcal{H}_1 . We then get

$$\sum_{m,n=1}^{\text{rank} E} \alpha_{mn} K_m^T K_n^* = 0,$$

which by condition Eq. (19) implies $\alpha_{mn} = 0$ for all m, n . Finally, by linear independence of $\{\sigma_a \otimes I, \sigma_a \otimes \sigma_b\}$, this also implies $\beta_a = 0 = \gamma_{ab}$ for all a and b . Conversely, one can write $|K_m\rangle\rangle\langle\langle K_n|$ as

$$|K_m\rangle\rangle\langle\langle K_n| = \frac{1}{d_1} I_1 \otimes K_m^T K_n^* + \Delta_{mn}, \quad (21)$$

where $\text{Tr}_1[\Delta_{mn}] = 0$ for all m, n . This implies that Δ_{mn} belongs to the span of $\mathbb{D}_{(N)}$ for all m, n . Let us suppose that Choi's condition Eq. (19) is not satisfied. Then there exist nontrivial coefficients ζ_{mn} such that $\sum_{m,n} \zeta_{m,n} K_m^T K_n^* = 0$. If we then take β_a, γ_{ab} such that

$$\sum_{mn} \zeta_{mn} \Delta_{mn} = \sum_a \beta_a \sigma_a \otimes I + \sum_{ab} \gamma_{ab} \sigma_a \otimes \sigma_b, \quad (22)$$

we have

$$\sum_{mn} \zeta_{mn} |K_m\rangle\rangle\langle\langle K_n| - \sum_a \beta_a \sigma_a \otimes I - \sum_{ab} \gamma_{ab} \sigma_a \otimes \sigma_b = 0, \quad (23)$$

in contradiction with Eq. (20).

VI. EXTREMALITY OF QUANTUM INSTRUMENTS

In contrast to a channel ($N = 1$, $M = 1$), which is specified by its Choi-Jamiolkowski operator, an instrument ($N = 1$, $M \geq 1$) is characterized by a collection of Choi-Jamiolkowski operators $\{N_i\}_{i=1}^M \subseteq \mathcal{L}(\mathcal{H}_1 \otimes \mathcal{H}_0)$, which sum up to Choi-Jamiolkowski operator of some channel R . The set $\mathbb{D}_{(N)} = \{\sigma_a \otimes I, \sigma_a \otimes \sigma_b\}$ from Theorem 1 is the same as for channels, because it depends only

on N , the number of teeth of GQI, but not on M the number of outcomes of the instrument. We can take the spectral decompositions of all the Choi-Jamiolkowski operators of the instrument $N_i = \sum_m |K_m^{(i)}\rangle\rangle\langle\langle K_m^{(i)}|$ and we can write the necessary and sufficient condition of extremality as follows.

Corollary 3: Instrument $\{N_i\}_{i=1}^M \subseteq \mathcal{L}(\mathcal{H}_1 \otimes \mathcal{H}_0)$ is extremal if and only if equation

$$\sum_{i,m,n} \alpha_{mn}^i |K_m^{(i)}\rangle\rangle\langle\langle K_n^{(i)}| + \sum_a \beta_a \sigma_a \otimes I + \sum_{a,b} \gamma_{ab} \sigma_a \otimes \sigma_b = 0 \quad (24)$$

cannot be satisfied for non-trivial coefficients $\alpha_{mn}^i, \beta_a, \gamma_{ab}$.

Counting the terms in Eq. (24) that have to be linearly independent elements of $\mathcal{L}(\mathcal{H}_1 \otimes \mathcal{H}_0)$, we can obtain a simple restriction on the ranks of the elements of the extremal instrument.

Corollary 4: An extremal instrument $\{N_i\}_{i=1}^M \subseteq \mathcal{L}(\mathcal{H}_1 \otimes \mathcal{H}_0)$ satisfies the following inequality

$$\sum_i r_i^2 \leq (d_0)^2, \quad (25)$$

where r_i denotes the rank of N_i and $d_0 = \dim \mathcal{H}_0$.

We will now prove a theorem that provides an equivalent, but more practical, extremality condition for quantum instruments.

Theorem 5: *An instrument $\{\mathcal{N}_i\}_{i=1}^M$ with Choi-Jamiolkowski operators $\{N_i = \sum_m |K_m^{(i)}\rangle\rangle\langle\langle K_m^{(i)}|\}_{i=1}^M$ is extremal if and only if the operators $\{K_m^{(i)\dagger} K_n^{(i)}\}$ are linearly independent.*

Proof: Suppose that the operators $\{K_m^{(i)\dagger} K_n^{(i)}\}$ are linearly independent. Then if Eq. (24) is satisfied, also its partial trace over space \mathcal{H}_1 is satisfied, namely

$$\sum_{i,m,n} \alpha_{mn}^i K_m^{(i)T} K_n^{(i)*} = 0, \quad (26)$$

which implies $\alpha_{mn}^i = 0$ for all i, m, n and consequently also $\beta_a = 0$ for all a and $\gamma_{ab} = 0$ for all a, b . Conversely, consider the extremality condition in Eq. (24) along with the following generalization of Eq. (21)

$$|K_m^{(i)}\rangle\rangle\langle\langle K_n^{(i)}| = \frac{1}{d_1} I_1 \otimes K_m^{(i)T} K_n^{(i)*} + \Delta_{mn}^{(i)}, \quad (27)$$

where the operators $\Delta_{mn}^{(i)}$ belong to the span of $\mathbb{D}_{(N)}$. If the operators $\{K_m^{(i)\dagger} K_n^{(i)}\}$ are not linearly independent, then there are non-trivial coefficients $\zeta_{mn}^{(i)}$ such that $\sum_{i,m,n} \zeta_{mn}^{(i)} K_m^{(i)T} K_n^{(i)*} = 0$. Then, taking β_a and γ_{ab} such that

$$\sum_{i,m,n} \zeta_{mn}^{(i)} \Delta_{mn}^{(i)} = \sum_a \beta_a \sigma_a \otimes I + \sum_{ab} \gamma_{ab} \sigma_a \otimes \sigma_b, \quad (28)$$

we have

$$\sum_{i,m,n} \zeta_{mn}^{(i)} |K_m^{(i)}\rangle\rangle\langle\langle K_n^{(i)}| - \sum_a \beta_a \sigma_a \otimes I - \sum_{ab} \gamma_{ab} \sigma_a \otimes \sigma_b = 0, \quad (29)$$

in contradiction with Eq. (24).

A. Extremality of Von Neuman-Lüders instruments

Let us now consider instruments of the following type

$$\mathcal{N}_i(\rho) = \sqrt{P_i} \rho \sqrt{P_i}, \quad (30)$$

where P_i is a POVM. Then, by Theorem 5, the instrument is extremal if and only if the POVM $\{P_i\}_{i=1}^M$ is linearly independent. Indeed, the set $\{K_m^{(i)\dagger} K_n^{(i)}\}$ in this case is provided precisely by $\{P_i\}_{i=1}^M$.

In particular, von Neuman-Lüders instruments are extremal. Indeed, every such instrument $\{N_i\}_{i=1}^d$ is of the form of Eq. (30) with $P_i = \Pi_i$ where $\Pi_i \Pi_j = \delta_{ij} \Pi_i$. Using the last constraint it is easy to prove that if $X = \sum_i \alpha_i \Pi_i = 0$ then $\Pi_j X = \alpha_j \Pi_j = 0$ and consequently $\alpha_j = 0$.

Since there exist POVMs that are not extremal, but have linearly independent elements, one can easily construct examples of extremal instruments, with non-extremal POVMs. For example, this is the case with $d = 2$ and $P_1 = 1/2|0\rangle\langle 0|$, $P_2 = 1/2|0\rangle\langle 0| + |1\rangle\langle 1|$.

VII. CONCLUSIONS

The aim of this paper was to characterize the extremal points of the set of GQIs. Our main result is represented by Theorem 1, which links extremality of the considered GQI with linear independence of a set of operators. An important special case of GQIs are Quantum testers. For quantum 1-testers we derived necessary and sufficient criterion of extremality that differs from the application of general Theorem 1 and can be tested more easily. As a consequence of the criterion, we obtained a bound (13) on the ranks of elements of the extremal 1-tester. We showed that the subsets of extremal 1-testers with a fixed normalization are isomorphic if they have the same rank of the normalization. This implies that to classify all extremal 1-testers it suffices to study extremal 1-testers with a completely mixed normalization ($I_2 \otimes \rho_1 = \frac{1}{d_1} I_{21}$). We completely characterized qubit 1-testers with 1 and 2 outcomes and provided techniques to construct extremal qubit testers with up to 13 outcomes, which is the maximal number allowed by the bound (14).

In Sec. V we apply our extremality condition from Theorem 1 to channels. The resulting condition is different from the well known criterion of Choi,¹⁶ even though we prove it to be equivalent. The Sec. VI presents the first characterization of the extremality of instruments. In particular, we show that instruments of the type defined in Eq. (30) for POVMs $\{P_i\}_{i=1}^M$ with linearly independent elements are extremal quantum instruments.

More generally, any quantum instrument determines not only a POVM, when the quantum output is ignored, but also a quantum channel, when the classical outcome is ignored. A natural question is then what combinations of extremality can exist when we consider an instrument along with the POVM and channel it defines. In Appendix C we present examples of instruments for seven out of the eight possibilities. The question whether non-extremal instruments exist, such that they determine extremal POVMs and extremal channels is left as an open problem.

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APPENDIX A: PARAMETRIZATION OF THE SET OF DETERMINISTIC QUANTUM COMBS

Suppose we want to choose such parametrization of the set of hermitian operators in which the subset of deterministic quantum combs would simply correspond to positive operators that have some parameters fixed (e.g., to zero). Let us consider a quantum N -combs $R \in \mathcal{L}(\mathcal{H}_{2N-1} \otimes \dots \otimes \mathcal{H}_0)$. For each of the Hilbert spaces \mathcal{H}_k $k \in \{0, \dots, 2N-1\}$ we choose a basis of hermitian operators on \mathcal{H}_k $\{E_a^{(k)}\}_{a=1}^{d_k^2}$ such that $E_1^{(k)} = I$ and all the other elements have zero trace. Taking the tensor product of the basis elements for all the Hilbert spaces \mathcal{H}_k we obtain a basis of hermitian operators $\{E_{a_{2N-1}}^{(2N-1)} \otimes \dots \otimes E_{a_0}^{(0)}\}$ on $\mathcal{H}_{2N-1} \otimes \dots \otimes \mathcal{H}_0$.

Let us now use this basis to illustrate the normalization cascade requirements on the quantum 1-combs, i.e., Choi operators of quantum channels. In this case a quantum channel mapping from $\mathcal{L}(\mathcal{H}_0)$ to $\mathcal{L}(\mathcal{H}_1)$ is represented via Choi-Jamiolkowski isomorphism by a positive operator $R \in$

$\mathcal{L}(\mathcal{H}_1 \otimes \mathcal{H}_0)$, which has to fulfil equation $\text{Tr}_1 R = I_0$. Using our basis arbitrary R can be written as

$$\begin{aligned} R &= \sum_{a_1=1}^{d_1^2} \sum_{a_0=1}^{d_0^2} c_{a_1 a_0} E_{a_1}^{(1)} \otimes E_{a_0}^{(0)} \\ &= c_{11} I_1 \otimes I_0 + I_1 \otimes \sum_{a_0=2}^{d_0^2} c_{1 a_0} E_{a_0}^{(0)} + \\ &\quad + \sum_{a_1=2}^{d_1^2} E_{a_1}^{(1)} \otimes \sum_{a_0=1}^{d_0^2} c_{a_1 a_0} E_{a_0}^{(0)} \end{aligned} \quad (\text{A1})$$

Let us now look how the three terms of the RHS of Eq. (A1) contribute to $\text{Tr}_1(R)$. The first two terms do contribute, whereas the remaining one does not. The requirement of $\text{Tr}_1(R) = I_0$ translates into the following equations $c_{11} = \frac{1}{d_1}$, $c_{1i} = 0 \forall i = 2, \dots, d_0$ for parameters $c_{a_1 a_0}$. Thus, each quantum 1-comb (Choi operator of a channel) can be written as

$$R = \frac{1}{d_1} I_{10} + \sum_i E_i^{(1)} \otimes A_i, \quad (\text{A2})$$

where A_i are arbitrary hermitian operators on \mathcal{H}_0 . Previous statements can be easily generalized to the case of general quantum combs. We shall first illustrate the relation of expansions for $R^{(n)}$ and $R^{(n-1)}$ and then write the expansion of general quantum N -comb. In our basis $R^{(n)}$ can be written as

$$\begin{aligned} R^{(n)} &= \sum_{a_{2n-1}, a_{2n-2}} E_{a_{2n-1}}^{(2n-1)} \otimes E_{a_{2n-2}}^{(2n-2)} \otimes L^{a_{2n-1}, a_{2n-2}}, \\ &= I_{2n-1, 2n-2} \otimes L^{1,1} + I_{2n-1} \otimes \sum_{j=2}^{d_{2n-2}^2} E_j^{(2n-2)} \otimes L^{1,j} \\ &\quad + \sum_{i=2}^{d_{2n-1}^2} E_i^{(2n-1)} \otimes \sum_{j=1}^{d_{2n-2}^2} E_j^{(2n-2)} \otimes L^{i,j} \end{aligned} \quad (\text{A3})$$

where

$$L^{a_{2n-1}, a_{2n-2}} = \sum_{a_{2n-3} \dots a_0} c_{a_{2n-1} \dots a_0} E_{a_{2n-3}}^{(2n-3)} \otimes \dots \otimes E_{a_0}^{(0)}$$

and we expanded the two sums in the same way as in (A1). The normalization cascade (1) requires²⁶ that

$$L^{1,1} = \frac{1}{d_{2n-1}} R_{(n-1)}, \quad L^{1,j} = 0 \forall j \quad (\text{A4})$$

and the operator $\sum_{j=1}^{d_{2n-2}^2} E_j^{(2n-2)} \otimes L^{i,j}$ can be an arbitrary operator on $\mathcal{H}_{2n-2} \otimes \dots \otimes \mathcal{H}_0$. As a result

$$R_{(n)} = \frac{1}{d_{2n-1}} I_{2n-1, 2n-2} \otimes R_{(n-1)} + \sum_{i=2}^{d_{2n-1}^2} E_i^{(2n-1)} \otimes B_i, \quad (\text{A5})$$

where $B_i \in \mathcal{L}(\mathcal{H}_{2n-2} \otimes \cdots \otimes \mathcal{H}_0)$. Using the above relation recursively we can write the parametrization of the general deterministic N -comb as

$$\begin{aligned}
 R_{(N)} &= \frac{1}{d_{2N-1}d_{2N-3} \dots d_1} I_{2N-1, \dots, 0} \\
 &+ \sum_{i=2} E_i^{(2N-1)} \otimes B_i^{(2N-2)} \\
 &+ I_{2N-1, 2N-2} \otimes \sum_{i=2} E_i^{(2N-3)} \otimes B_i^{(2N-4)} + \dots \\
 &+ I_{2N-1, \dots, 2} \otimes \sum_{i=2} E_i^{(1)} \otimes B_i^{(0)}, \tag{A6}
 \end{aligned}$$

where $B_i^{(k)}$ are arbitrary hermitian operators acting on $\mathcal{H}_k \otimes \mathcal{H}_{k-1} \otimes \cdots \otimes \mathcal{H}_0$. Let us denote the basis for hermitian operators $B_i^{(k)}$ as $\{F_j^{(k)}\}_{j=1}^{d_k \dots d_0^2}$. Consequently, the basis used for the variable part (i.e., all terms in (A6) except the first) of the quantum comb is $\{E_i^{(2N-1)} \otimes F_j^{(2N-2)}, I_{2N-1, 2N-2} \otimes E_i^{(2N-3)} \otimes F_j^{(2N-4)}, \dots, I_{2N-1, \dots, 2} \otimes E_i^{(1)} \otimes F_j^{(0)}\}$ and we denote it as $\mathbb{D}_{(N)}$. The operator basis $\mathbb{R}_{(N)}$ sufficient to expand arbitrary deterministic comb is then formed by $\{I_{2N-1, \dots, 0}\} \cup \mathbb{D}_{(N)}$.

APPENDIX B: TWO OUTCOME QUBIT 1-TESTERS

Suppose we have a two-outcome qubit 1-tester $\{T_i = \frac{1}{2} P_i\}_{i=1}^2$ with normalization $I_2 \otimes \frac{1}{2} I_1$ and P_i being orthogonal projectors. Equivalently to Theorem 2 we can say that the two outcome 1-tester is extremal if and only if $\mathcal{V}_\sigma \cap \mathcal{V}_T = 0$, where $\mathcal{V}_\sigma = \text{Span}\{I \otimes \sigma_k\}_{k=x,y,z}$ and \mathcal{V}_T is the direct sum of the two subspaces of hermitian operators with support in P_1 and in P_2 , respectively. The non existence of the intersection of \mathcal{V}_σ and \mathcal{V}_T can be stated also as the impossibility to fulfill the following equation

$$\sum_{i=1}^4 \lambda_i |\phi_i\rangle\langle\phi_i| = I \otimes (n_x \sigma_x + n_y \sigma_y + n_z \sigma_z), \tag{B1}$$

where the left hand side of Eq. (B1) represents a generic element in \mathcal{V}_T and the right hand side a generic element of \mathcal{V}_σ . The set of $\{|\phi_i\rangle\}_{i=1}^4$ forms an orthonormal basis of vectors belonging to $\text{Supp}(P_1) \cup \text{Supp}(P_2)$, and without loss of generality we can take $n_x^2 + n_y^2 + n_z^2 = 1$. This guarantees that the RHS has the spectral decomposition of the following form:

$$I \otimes |v\rangle\langle v| - I \otimes |v^\perp\rangle\langle v^\perp|,$$

with two +1 eigenvalues and two -1 eigenvalues and vector $|v\rangle$ that can be arbitrary thanks to freedom in n_x, n_y, n_z . Moreover projectors P_i can be written as $P_1 = \sum_{i=1}^{r_1} |\phi_i\rangle\langle\phi_i|$ and $P_2 = \sum_{i=r_1+1}^4 |\phi_i\rangle\langle\phi_i|$. Let us now investigate the circumstances under which the equation can be fulfilled, i.e., the tester is not extremal.

1. Case $(r_1, r_2) = (1, 3)$

This type of tester must have the form $\{T_1 = \frac{1}{2} |\phi_1\rangle\langle\phi_1|, T_2 = \frac{1}{2} \sum_{i=2}^4 |\phi_i\rangle\langle\phi_i|\}$. The LHS of Eq. (B1) must have the same eigenvalues as the RHS. Without loss of generality we can assume $\lambda_1 = \lambda_2 = -\lambda_3 = -\lambda_4$, because we can suitably relabel $|\phi_2\rangle, |\phi_3\rangle, |\phi_4\rangle$. Hence, we have

$$\begin{aligned}
 |\phi_1\rangle\langle\phi_1| + |\phi_2\rangle\langle\phi_2| &= I \otimes |e\rangle\langle e|, \\
 |\phi_3\rangle\langle\phi_3| + |\phi_4\rangle\langle\phi_4| &= I \otimes |e^\perp\rangle\langle e^\perp|, \tag{B2}
 \end{aligned}$$

where $e = v$ or $e = v^\perp$ depending on $\lambda_1 = +1$ or $\lambda_1 = -1$, respectively. In both cases Eq. (B2) implies that the qubit 1-tester of the form $T_1 = \frac{1}{2}|\phi_1\rangle\langle\phi_1|$ $T_2 = \frac{1}{2}(I \otimes I - |\phi_1\rangle\langle\phi_1|)$ is not extremal if and only if $|\phi_1\rangle = |f\rangle \otimes |e\rangle$ is a product vector.

2. Case $(r_1, r_2) = (2, 2)$

In this case the tester has the form $\{T_1 = \frac{1}{2}P_1 = \frac{1}{2}(|\phi_1\rangle\langle\phi_1| + |\phi_2\rangle\langle\phi_2|)$, $T_2 = \frac{1}{2}P_2 = \frac{1}{2}(|\phi_3\rangle\langle\phi_3| + |\phi_4\rangle\langle\phi_4|)$. In order to fulfill the equation (B1) two λ_i 's must be equal to $+1$ and two to -1 . Thus, λ_1, λ_2 have either same signs or different signs. If $\lambda_1 = \lambda_2 = \pm 1$ then the equation (B1) can be fulfilled if and only if $P_1 = I \otimes |e\rangle\langle e|$, where $e = v$ for $\lambda_{1,2} = 1$ or $e = v^\perp$ for $\lambda_{1,2} = -1$. If $\lambda_1 = -\lambda_2$ then we can assume without loss of generality that $|\phi_3\rangle, |\phi_4\rangle$ are labeled so that $\lambda_1 = \lambda_3$. We have

$$|\phi_1\rangle\langle\phi_1| + |\phi_3\rangle\langle\phi_3| = I \otimes |e\rangle\langle e|, \quad (\text{B3})$$

where $e = v$ or $e = v^\perp$ depending on $\lambda_1 = 1$ or $\lambda_1 = -1$, respectively. This may hold only if $|\phi_1\rangle = |f\rangle \otimes |e\rangle$ and $|\phi_3\rangle = |f^\perp\rangle \otimes |e\rangle$ for some vector $|f\rangle \in \mathcal{H}_2$. Due to equation (B3) $|\phi_2\rangle\langle\phi_2| + |\phi_4\rangle\langle\phi_4| = I \otimes |e^\perp\rangle\langle e^\perp|$ and $|\phi_2\rangle = |h\rangle \otimes |e^\perp\rangle$, $|\phi_4\rangle = |h^\perp\rangle \otimes |e^\perp\rangle$ for some $|h\rangle \in \mathcal{H}_2$. Thus, if $\lambda_1 = -\lambda_2$ the tester is not extremal if and only if

$$P_1 = |f\rangle\langle f| \otimes |e\rangle\langle e| + |h\rangle\langle h| \otimes |e^\perp\rangle\langle e^\perp|$$

$$P_2 = |f^\perp\rangle\langle f^\perp| \otimes |e\rangle\langle e| + |h^\perp\rangle\langle h^\perp| \otimes |e^\perp\rangle\langle e^\perp|.$$

for some $|e\rangle \in \mathcal{H}_1$, $|f\rangle, |h\rangle \in \mathcal{H}_2$.²⁷ The form of projectors P_1, P_2 can be equivalently stated as the existence of a product vector $|f\rangle \otimes |e\rangle$ in the support of P_1 such that $|f^\perp\rangle \otimes |e\rangle$ belongs to the support of P_2 . From our derivation it should be clear that if $P_1 \neq I \otimes |e\rangle\langle e|$ for any $|e\rangle \in \mathcal{H}_1$ and P_1 does not have the above mentioned form then $\{T_1, T_2\}$ is an extremal qubit two-outcome tester.

APPENDIX C: EXTREMALITY OF AN INSTRUMENT AND THE POVM AND THE CHANNEL DERIVED FROM IT

The present Appendix addresses the question about possible combinations of extremality of an instrument and the POVM and channel derived from it. We show feasibility of seven out of the eight possible combinations, by providing an example for each of them. In the following table, we enumerate the possible combinations, and we define them by writing $+$ if the object in the corresponding column (channel, POVM, instrument) is extremal, and $-$ otherwise.

Combination	Instrument	Channel	POVM
1	+	+	+
2	+	+	-
3	+	-	+
4	+	-	-
5	-	+	+
6	-	+	-
7	-	-	+
8	-	-	-

The existence of an instrument corresponding to combination number 5 is left as an open problem. Here is a list of examples for each of the remaining combinations.

Combination 1: The identical transformation is the most simple example of this kind. Also constant mapping to a fixed pure state has the desired properties of extremality.

Combination 2: Consider an instrument with two outcomes mapping a single qubit into two-qubits. First, we define

$$\begin{aligned} P_0 &:= \frac{1}{3}|0\rangle\langle 0| + \frac{2}{3}|1\rangle\langle 1| \\ P_1 &:= \frac{2}{3}|0\rangle\langle 0| + \frac{1}{3}|1\rangle\langle 1| \\ W &:= |0\rangle\langle 1| - |1\rangle\langle 0| \end{aligned}$$

and we define the Kraus operators of the instrument as follows:

$$\begin{aligned} M_0 &:= \sqrt{P_0} \otimes \frac{1}{\sqrt{2}}(|0\rangle + |1\rangle) \\ M_1 &:= \frac{1}{\sqrt{2}}\sqrt{P_1} \otimes |0\rangle + \frac{1}{\sqrt{2}}W\sqrt{P_1} \otimes |1\rangle, \end{aligned}$$

where for each outcome $i = 0, 1$ we have only a single Kraus operator. One can easily verify, that the induced POVM $\{P_0, P_1\}$ is not extremal, but linear independence of its elements guarantees extremality of the instrument. In order to check extremality of the induced channel one needs to take the minimal Kraus representation and check Choi's linear independence condition.

Combination 3: The Lüders instrument of a Von Neumann measurement is an extremal instrument, which induces extremal POVM and a non extremal channel.

Combination 4: This desired type of instrument can be constructed as in Eq. (30) with a POVM $\{P_i\}_{i=1}^M$, whose elements are linearly independent and commute. As a simple example one can take the qubit POVM $\{P_0, P_1\}$ defined in Combination 2.

Combination 6: This type of instrument can be constructed as follows. One takes an extremal channel, whose minimal dilation has N (more than one) Kraus operators K_i . Using these operators we define two instruments with N outcomes differing only in the choice of Kraus operators that correspond to each outcome (e.g., $M_i^{(1)} = K_i, M_i^{(2)} = K_{\sigma(i)}$, where σ is a permutation). Taking a convex combination of the two instruments provides the desired example.

Combination 7: This type of instrument can be constructed as a convex combination of two instruments, which induce the same POVM, but different channels. One takes for example the instrument $\{\mathcal{N}_i\}_{i=1}^M$ as in Eq. (30) for an extremal POVM $\{P_i\}_{i=1}^M$, and mixes it with the same instrument, which in addition applies an unitary channel U on the quantum output. Obviously, the induced channel differs, while the induced POVM remains the same.

Combination 8: For the construction of this example it is sufficient to take a convex combination of two instruments, which induce different POVMs, and different channels.

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²⁶After performing the trace, we are comparing the operators expanded in the same basis.

²⁷Notice that for $|f\rangle = |h\rangle$ this form gives $P_1 = |f\rangle\langle f| \otimes I$. Moreover, the freedom in $|e\rangle$ comes from n_x, n_y, n_z and the freedom in $|f\rangle, |h\rangle$ relates to splitting of unity.