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### Principles for Quantum Theory

Giacomo Mauro D'Ariano Università degli Studi di Pavia

QMS: decoherence and empirical estimates

June 29 - July 1st 2015

- The experience in Quantum Information has led us to look at Quantum Theory (QT) under a completely new angle
- QT is a theory of information



Giulio Chiribella



Paolo Perinotti



### Informational derivation of quantum theory

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QUIT Group, Dipartimento di Fisica "A. Volta" and INFN Sezione di Pavia, via Bassi 6, I-27100 Pavia, Italy (Received 29 November 2010; published 11 July 2011)

We derive quantum theory from purely informational principles. Five elementary axioms—causality, perfect distinguishability, ideal compression, local distinguishability, and pure conditioning—define a broad class of theories of information processing that can be regarded as standard. One postulate—purification—singles out quantum theory within this class.

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### Principles for Quantum Theory

- P1. Causality
- P2. Local discriminability
- P3. Purification \*
- P4. Atomicity of composition
- P5. Perfect distinguishability
- P6. Lossless Compressibility

Book from CUP soon

- The experience in Quantum Information has led us to look at Quantum Theory (QT) under a completely new angle
- QT is a theory of information



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### Quantum purity: How the big picture banishes weirdness

08 April 2015 by **Anil Ananthaswamy**Magazine issue 3016. **Subscribe and save**For similar stories, visit the **Quantum World** Topic Guide



(Image: Julien Pacaud)

WE HAVE become accustomed to the universe blowing our minds – perhaps too accustomed. Quantum weirdness – things like particles being in two places at once, or appearing to share a telepathic link – has been baffling us for more than a century now. The physicist Richard Feynman once said that nobody really understands the quantum world. Or as others have put it: if you think you understand it, then you definitely don't. So it is tempting to throw up our hands and say human brains can never grasp it.

But maybe we shouldn't be so defeatist. Isn't it just possible that we simply haven't yet got to the bottom of how quantum mechanics really works? That's what Giacomo Mauro D'Ariano of the University of Pavia in Italy, and his colleagues Giulio Chiribella and Paolo Perinotti think – and they have been doing something ...

The framework

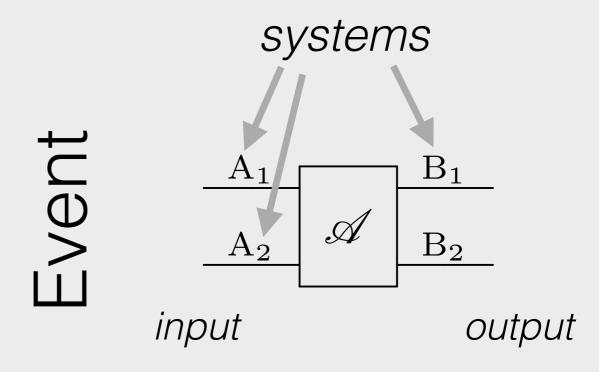
Logic c Probability c OPT

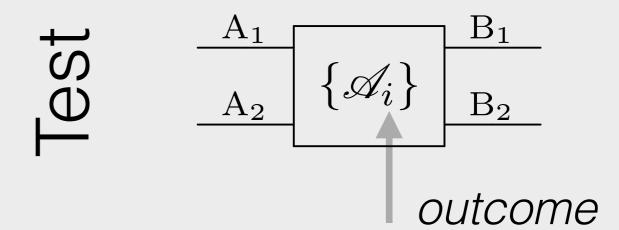
joint probabilities + connectivity

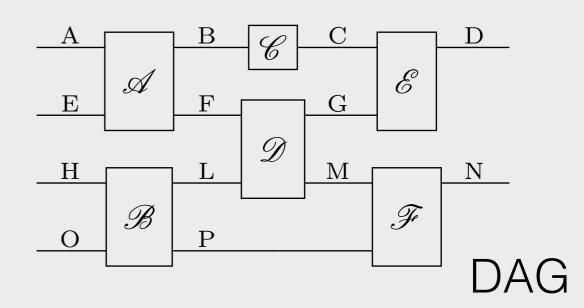
$$p(i, j, k, ... | \text{circuit})$$

Marginal probability

$$\sum_{i,k,...} p(i,j,k,...|\text{circuit}) = p(j|\text{circuit})$$







The framework

Logic c Probability c OPT

joint probabilities + connectivity

$$p(i, j, k, ... | \text{circuit})$$

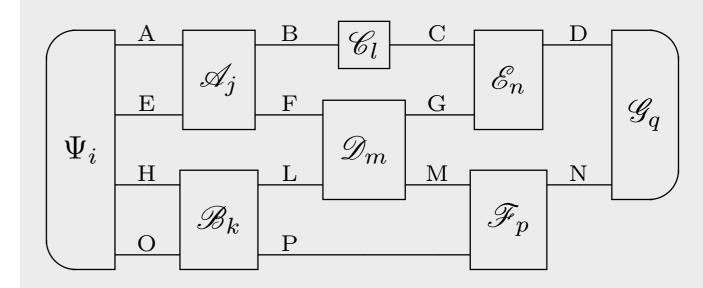
Notice: the probability of a "preparation" generally depends on the circuit at its output.

$$\begin{array}{c|c}
\hline
\rho_i
\end{array} := 
\begin{array}{c|c}
\hline
I
\end{array}$$

preparation

$$A a_j := A a_j$$

observation



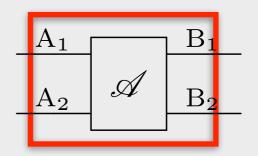
The framework

Logic c Probability c OPT

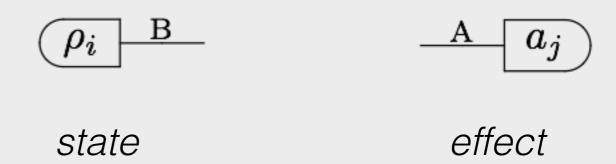
joint probabilities + connectivity

Probabilistic equivalence classes

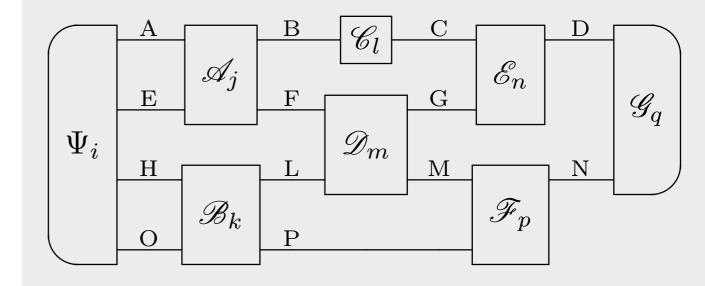
Notice: the probability of a transformation generally depends on the circuit at its output!!



transformation



p(i, j, k, l, m, n, p, q | circuit)

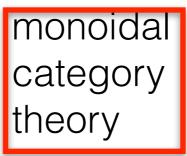


The framework

Logic c Probability c OPT

joint probabilities + connectivity

Probabilistic equivalence classes



Multiplication of closed circuits

$$\begin{array}{c|c}
\hline
\rho_{i_1} & A & a_{i_2} \\
\hline
\sigma_{j_1} & B & b_{j_2}
\end{array} = 
\begin{array}{c|c}
\hline
\rho_{i_1} & A & a_{i_2} \\
\hline
\sigma_{j_1} & B & b_{j_2}
\end{array}$$

$$= p(i_1, i_2) q(j_1, j_2)$$

Sequential composition (associative)

Identity test

Parallel composition (associative)

$$AB = BA$$

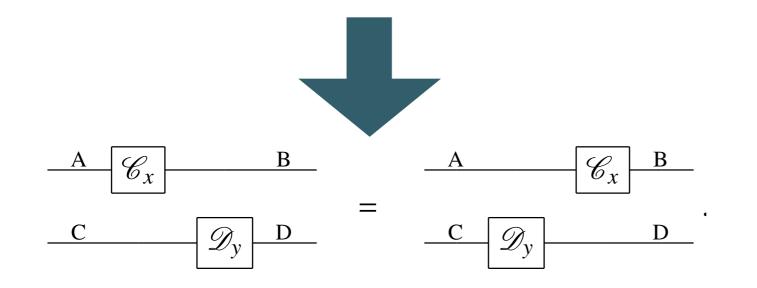
$$\{\rho_x\}_{x \in X} \xrightarrow{B} := \frac{I}{\{\rho_x\}_{x \in X}} \xrightarrow{B}$$

$$AI = IA = A$$

$$A(BC) = (AB)C \qquad \xrightarrow{A} \{a_y\}_{y \in Y} := \xrightarrow{A} \{a_y\}_{y \in Y}$$

Sequential and parallel compositions commute

$$(\mathscr{A}\otimes\mathscr{D})\circ(\mathscr{C}\otimes\mathscr{B})=(\mathscr{A}\circ\mathscr{C})\otimes(\mathscr{D}\circ\mathscr{B})$$



<del>-</del>

wire-stretching

(foliations)

The framework

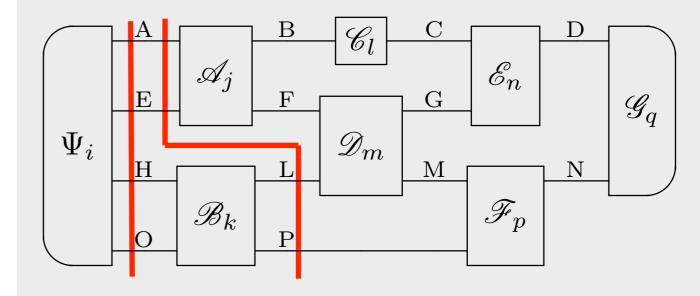
Logic c Probability c OPT

joint probabilities + connectivity

p(i, j, k, ... | circuit)

Maximal set of independent systems = "leaf"

$$p(i, j, k, l, m, n, p, q | \text{circuit})$$



The framework

Logic c Probability c OPT

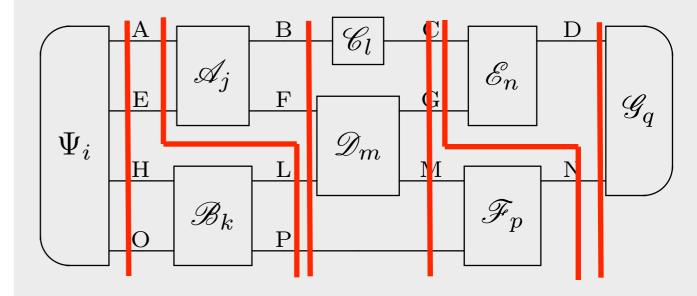
joint probabilities + connectivity

p(i, j, k, ... | circuit)

Maximal set of independent systems = "leaf"

Foliation

p(i, j, k, l, m, n, p, q | circuit)



States are functionals for effects

States are separating for effects

Effects are functionals on states

Effects are separating for states

Embedding in real vector spaces

 $\mathsf{St}(\mathsf{A}),\;\mathsf{St}_1(\mathsf{A}),\;\mathsf{St}_\mathbb{R}(\mathsf{A})$ 

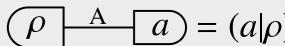
 $\mathsf{Eff}(\mathsf{A}),\;\mathsf{Eff}_1(\mathsf{A}),\;\mathsf{Eff}_\mathbb{R}(\mathsf{A})$ 

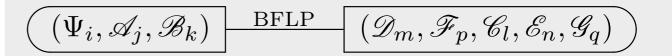
Dimension  $D_{
m A}$ 

$$\mathsf{Eff}_{\mathbb{R}}(\mathsf{A}) = \mathsf{St}_{\mathbb{R}}(\mathsf{A})^{\vee}$$
  $\mathsf{St}_{\mathbb{R}}(\mathsf{A}) = \mathsf{Eff}_{\mathbb{R}}(\mathsf{A})^{\vee}$ 

### Paring notation:

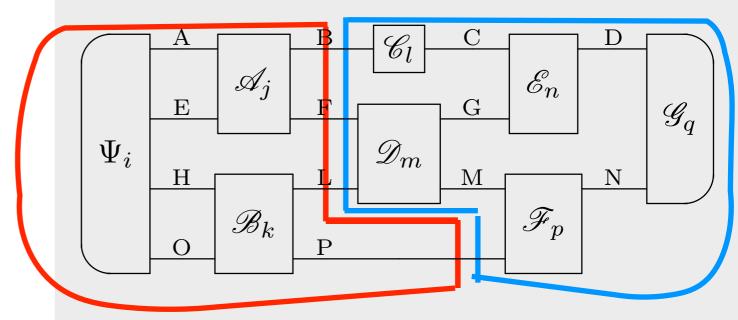
$$\rho \in St(A), a \in Eff(A), \quad \boxed{\rho}$$







p(i, j, k, l, m, n, p, q | circuit)



$$\{\mathcal{T}_i\}_{i\in\{\underbrace{i_1,i_2,\ldots,i_n,i_{n+1},i_{n+2},\ldots,\ldots}_{j_1}\}}$$

Coarse-graining



$$\{\hat{\mathcal{T}}_j\}_{j\in\{j_1,j_2,\ldots\}}$$

$$\hat{\mathscr{T}}_{S} = \sum_{i \in S} \mathscr{T}_{i}$$

Partial ordering

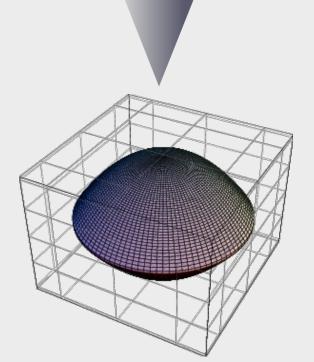
### Conditioned test (needs causality)

### Circuit multiplication: randomize tests

$$p_i \stackrel{\mathrm{A}}{\longrightarrow} \mathscr{C}^{(i)}_{j_i} \stackrel{\mathrm{B}}{\longrightarrow} := \underbrace{ \stackrel{\mathrm{A}}{\longrightarrow} \underbrace{ \mathscr{C}^{(i)}_{j_i} \stackrel{\mathrm{B}}{\longrightarrow} }_{I} }_{I}$$



Cone structure



Convex structure

State tomography

 $\{l_i\}_{i\in X}\subseteq \mathsf{Eff}(A)$  separating for states  $\longrightarrow$  span  $\mathsf{Eff}(A)$ 





$$\forall a \in \mathsf{Eff}(\mathsf{A}), \ a = \sum_{i \in \mathsf{X}} c_i(a) l_i \qquad c_i \in \mathsf{St}_{\mathbb{R}}(\mathsf{A}).$$
  $\{c_i\}_{i \in \mathsf{X}} \text{ is a dual set for } \{l_i\}_{i \in \mathsf{X}}$ 

$$ho\in \mathrm{St}_1(\mathrm{A})$$
 deterministic  $orall a\in \mathrm{Eff}_{\mathbb{R}}(\mathrm{A}),\; (a|
ho)=\sum_{i\in \mathrm{X}}c_i(a)(l_i|
ho)$  state-tomography



 $\{l_i\}_{i\in X}$  informationally complete for states

 $\{\rho_0, \rho_1\} \subseteq St(A)$  preparation test

$$\{a_0, a_1\}$$

observation test

success probability of discrimination

$$p_{\text{succ}} = (a_0|\rho_0) + (a_1|\rho_1)$$

$$= (a|\rho_0) + (a_1|\rho_1 - \rho_0)$$

$$= (a|\rho_1) + (a_0|\rho_0 - \rho_1)$$

$$= \frac{1}{2}[1 + (a_1 - a_0|\rho_1 - \rho_0)]$$

Metric

$$p_{\text{succ}}^{(\text{opt})} = \frac{1}{2}[1 + ||\rho_1 - \rho_0||]$$

$$\|\delta\| := \sup_{\{a_0,a_1\}} (a_0 - a_1 | \delta),$$

$$\|\delta\| = \sup_{a_0 \in \mathsf{Eff}(\mathsf{A})} (a_0|\delta) - \inf_{a_1 \in \mathsf{Eff}(\mathsf{A})} (a_1|\delta)$$

monotonicity

$$\mathscr{C} \in \mathsf{Transf}_1(A, B)$$

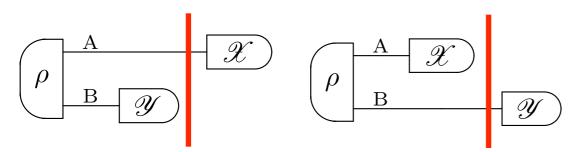
$$\|\mathscr{C}\delta\|_{\mathrm{B}} \leq \|\delta\|_{\mathrm{A}}$$

$$a := a_0 + a_1$$

- P1. Causality
- P2. Local discriminability
- P3. Purilication
- P4. Atomicity of composition
- P5. Perfect distinguishability
- P6. Lossless Compressibility

The probability of preparations is independent of the choice of observations





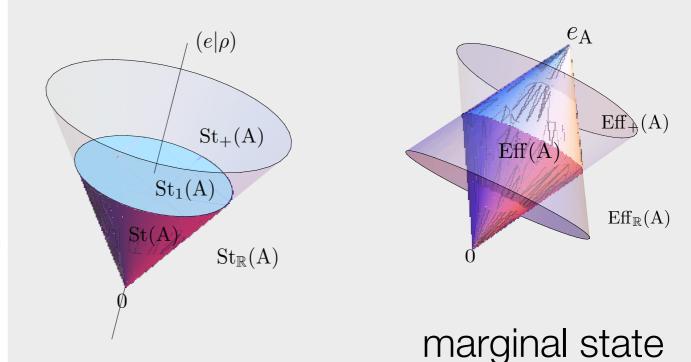


$$p(i, j|\mathcal{X}, \mathcal{Y}) := (a_j|\rho_i)$$



$$p(i|\mathscr{X},\mathscr{Y}) = p(i|\mathscr{X},\mathscr{Y}') = p(i|\mathscr{X})$$

Iff conditions: a) the deterministic effect is unique; b) states are "normalizable"



- P1. Causality
- P2. Local discriminability
- P3. Purification
- P4. Atomicity of composition
- P5. Perfect distinguishability
- P6. Lossless Compressibility

It is possible to discriminate any pair of states of composite systems using only local measurements.



Origin of the complex tensor product

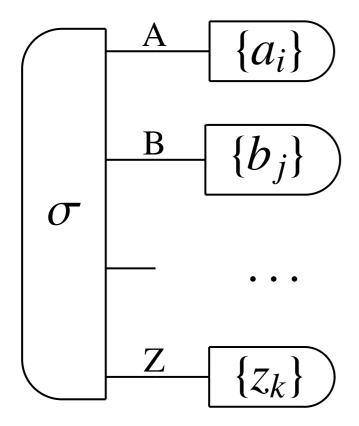


Local characterization of transformations

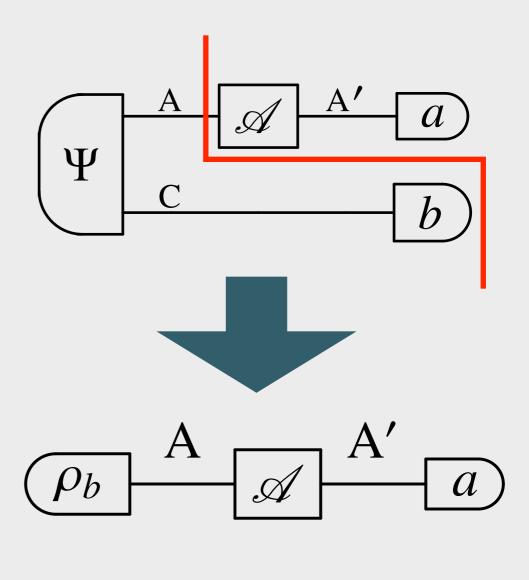
$$\begin{array}{c|c}
A & A' \\
\hline
\Psi & B \\
\hline
b \\
\hline
\end{array} = \begin{array}{c}
\rho_b & A' \\
\hline
A' & a
\end{array}$$



Local effects are separating for joint states



### Tomography

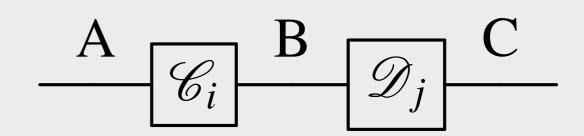


- P1. Causality
- P2. Local discriminability
- P3. Purification
- P4. Atomicity of composition
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The composition of two atomic transformations is atomic



Complete information can be accessed on a step-by-step basis

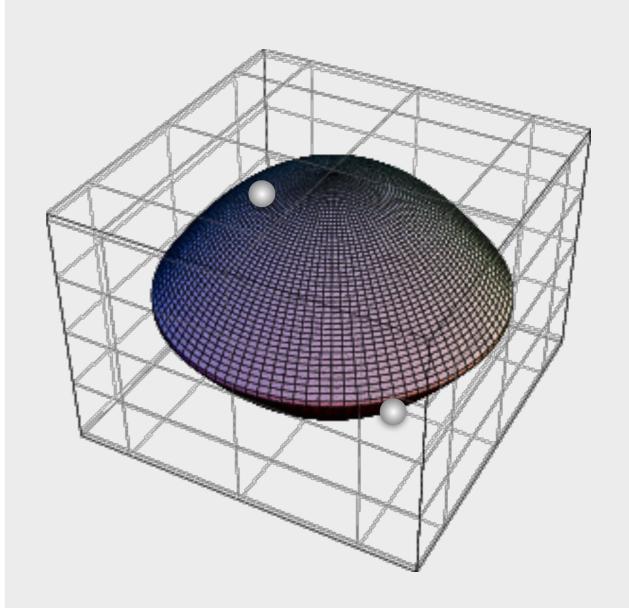


- P1. Causality
- P2. Local discriminability
- P3. Purification
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Every state that is not completely mixed (i.e. on the boundary of the convex) can be perfectly distinguished from some other state.



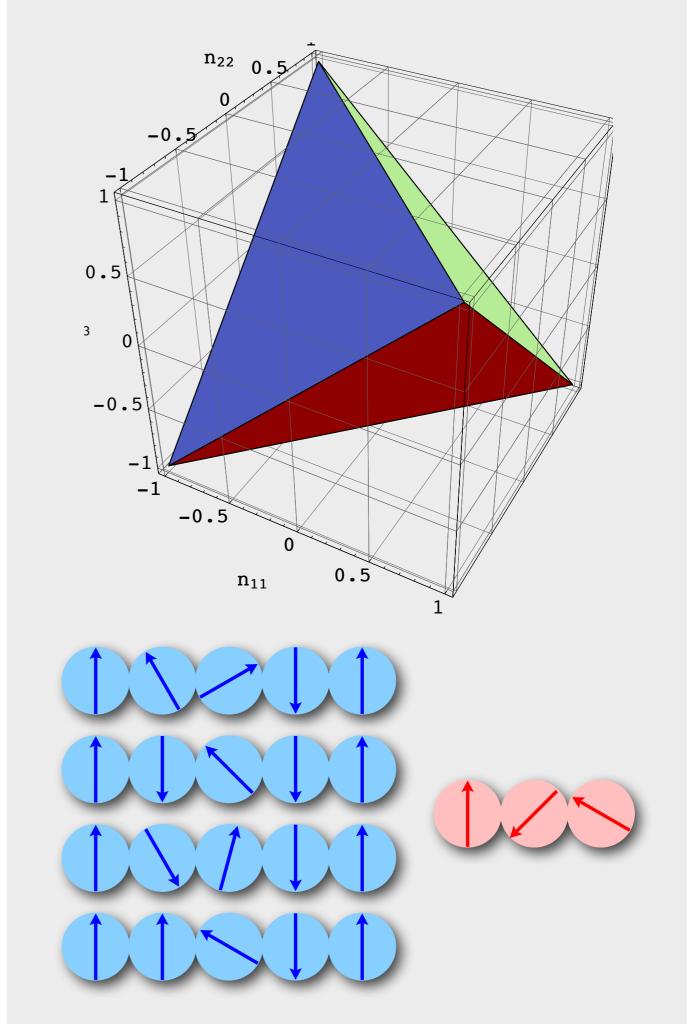
Falsifiability of the theory



- P1. Causality
- P2. Local discriminability
- P3. Purification
- P4. Atomicity of composition
- P5. Perfect distinguishability
- P6. Lossless Compressibility

For states that are not completely mixed there exists an ideal compression scheme

Any face of the convex set of states is the convex set of states of some other system



- P1. Causality
- P2. Local discriminability
- P3. Purification
- P4. Atomicity of composition
- P5. Perfect distinguishability
- P6. Lossless Compressibility

Every state has a purification. For fixed purifying system, every two purifications of the same state are connected by a reversible transformation on the purifying system

$$\rho A = \Psi B e$$

$$\Psi' B e$$

$$\Psi' B e$$

$$\Psi' B e$$

$$\Psi' B e$$

### Principles for

### Quantum Theory

- P1. Causality
- P2. Local discriminability
- P3. Purification
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- P6. Lossless Compressibility

Every state has a purification. For fixed purifying system, every two purifications of the same state are connected by a reversible transformation on the purifying system

### 1. Existence of entangled states:

the purification of a mixed state is an entangled state; the marginal of a pure entangled state is a mixed state;

2. Every two normalized pure states of the same system are connected by a reversible transformation

$$\boxed{\psi'} = \boxed{\psi} = \boxed{\mathcal{U}} = \boxed{\mathcal{U}}$$

3. **Steering:** Let  $\Psi$  purification of  $\rho$ . The for every ensemble decomposition  $\rho = \sum_{x} p_{x} \alpha_{x}$  there exists a measurement  $\{b_{x}\}$ , such that

4. Process tomography (faithful state):

5. No information without disturbance

- P1. Causality
- P2. Local discriminability
- P3. Purification
- P4. Atomicity of composition
- P5. Perfect distinguishability
- P6. Lossless Compressibility

Every state has a purification. For fixed purifying system, every two purifications of the same state are connected by a reversible transformation on the purifying system

### Consequences

### 6. Teleportation

### 7. Reversible dilation of "channels"

$$\begin{array}{c}
A & \mathcal{C} & A \\
\hline
A & \mathcal{C} & A
\end{array} = 
\begin{array}{c}
\eta & E & E & E \\
\hline
A & \mathcal{U} & A
\end{array}$$

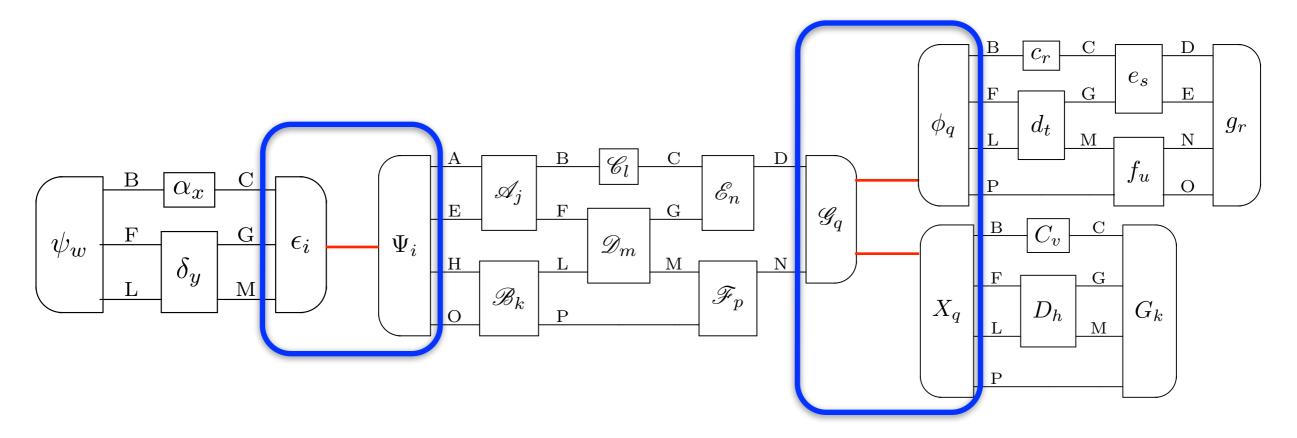
### 8. Reversible dilation of "instruments"

### 9. State-transformation cone isomorphism

10. Rev. transform. for a system make a compact Lie group

### On the von Neumann postulate





P1. Causality

P2. Local discriminability

P3. Purification

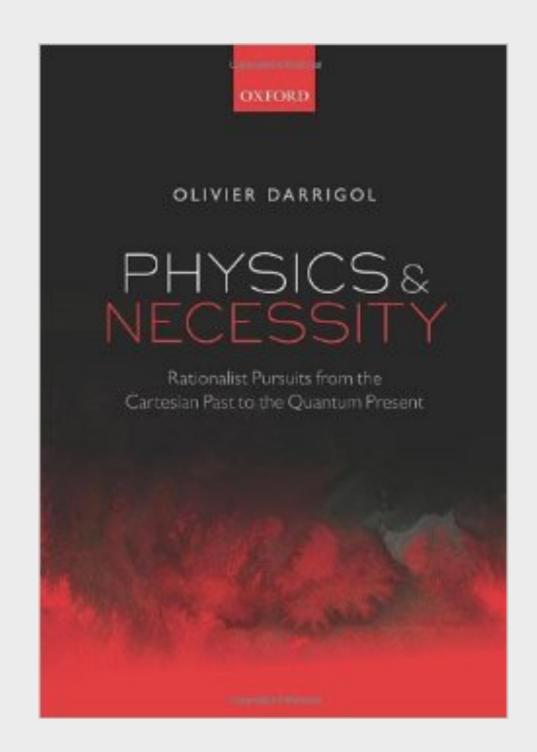
P4. Atomicity of composition

P5. Perfect distinguishability

P6. Lossless Compressibility

Epistemological principles
Are they necessary?

Fermionic quantum theory?



Thank you!